

Sea Ice-Albedo Feedback and Nonlinear Arctic Climate Change

Michael Winton

Geophysical Fluid Dynamics Laboratory/NOAA

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Corresponding author: Dr. Michael Winton, GFDL/NOAA, P.O. Box 308, Princeton University

Forrestal Campus, Princeton, NJ 08542. email: Michael.Winton@noaa.gov

Abstract

The potential for sea ice-albedo feedback to give rise to nonlinear climate change in the Arctic Ocean – defined as a nonlinear relationship between polar and global temperature change or, equivalently, a time-varying polar amplification – is explored in IPCC AR4 climate models. Five models supplying SRES A1B ensembles for the 21st century are examined and very linear relationships are found between polar and global temperatures (indicating linear Arctic Ocean climate change), and between polar temperature and albedo (the potential source of nonlinearity). Two of the climate models have Arctic Ocean simulations that become annually sea ice-free under the stronger CO₂ increase to quadrupling forcing. Both of these runs show increases in polar amplification at polar temperatures above -5°C and one exhibits heat budget changes that are consistent with the *small ice cap instability* of simple energy balance models. Both models show linear warming up to a polar temperature of -5°C, well above the disappearance of their September ice covers at about -9°C. Below -5°C, surface albedo decreases smoothly as reductions move, progressively, to earlier parts of the sunlit period. Atmospheric heat transport exerts a strong cooling effect during the transition to annually ice-free conditions. Specialized experiments with atmosphere and coupled models show that the main damping mechanism for sea ice region surface temperature is reduced upward heat flux through the adjacent ice-free oceans resulting in reduced atmospheric heat transport into the region.

I. Introduction

The speculation that Arctic climate has nonlinear behaviors associated with sea ice albedo feedback has deep roots in climatology (Brooks 1949, Donn and Ewing 1968). Energy balance models (EBMs) were used to study ice albedo effects starting in the late 1960s and one of the first uses of atmospheric global climate models (GCMs) was to explore the climatic impact of the Arctic sea ice cover (see review in Royer et al, 1990). Although climate models show that global temperature change is mainly linear in climate forcing over a broad range (Hansen et al, 2005), the nonlinear relationship between ice albedo and temperature may introduce local nonlinearity. Simple diffusive energy balance models, that represent this relationship with a step function, produce an abrupt disappearance of polar ice as the global climate gradually warms (North 1984). The phenomenon is known as the *small ice cap instability* (SICI) as it disallows polar ice caps smaller than a certain critical size related to heat diffusion and radiative damping parameters. Thorndike (1992) coupled an atmospheric energy balance model to a simple analytical model of sea ice in an ocean mixed layer, thereby simulating rather than parameterizing the albedo temperature relationship, and found that seasonally ice-free states were unstable. Under increased forcing, Thorndike's "toy" model transitions directly from annually ice-covered to annually ice-free states inducing a large and abrupt increase in surface temperature.

The Arctic sea ice cover has been in decline since the 1950s (Vinnikov et al, 1999). This decline is more pronounced in the summer and recent years have produced striking record minima (Stroeve et al, 2005). Some researchers have noted that nonlinear behaviors such as thresholds and tipping points may be associated with this decline

(Lindsay and Zhang 2005; Serreze and Francis 2006). The goal of this paper is to assess the potential for nonlinearity of Arctic climate change in the IPCC fourth assessment report (AR4) climate models. In section II we demonstrate potential nonlinearities in a simple model and develop a strategy for assessment. In section III we examine 21st century simulations for signs of nonlinearity. Section IV continues this search by examining two strongly forced experiments as they become annually ice-free. Section V shows that the nonlinear behavior of one of these experiments is similar to the EBM SICI. Section VI explores the stabilizing effect of ocean surface fluxes and atmospheric heat transport on the sea ice with special GCM experiments designed to illuminate the climate response to sea ice region changes. Section VIII summarizes and discusses the results.

II. Elementary Arctic climate dynamics

The potential for a nonlinear relationship between ice albedo and temperature to generate nonlinear climate change can be demonstrated with a very simple energy balance model. Consider the energy balance at the top of an isolated polar atmosphere:

$$A+BT=S[1-\alpha(T)] \quad (1)$$

The model represents a balance between absorbed shortwave, insolation (S) times a planetary albedo ($1-\alpha$), and parameterized outgoing longwave radiation with a linear dependence on surface temperature, T . The model is isolated in the sense that the atmospheric heat transport convergence is held fixed – bundled with the longwave intercept into A . The nonlinearity of the model comes from the nonlinear dependence of α on T . At very low mean temperatures, where snow never melts, albedo is insensitive to temperature. The same is true at high mean temperatures where there is no ice. Between

these flat sections, there is a drop from snow to seawater albedos. One might expect, based on the liquid/ice transition occurring at a fixed temperature, that this drop would resemble a cliff. However, the seasonal cycle and other variability allow sampling of various ice-cover states at any given long-term mean temperature, smoothing the relationship. For simplicity, let's take this smoothed section to be linear and call it the ramp. The slope of the ramp depends on the drop in albedo between its endpoints and the temperature range over which the drop is experienced. The drop in planetary albedo, the albedo above the atmosphere, will be less than the jump in surface albedo because only part of the insolation reaches and interacts with the surface. There may also be changes in atmospheric properties with temperature that impact the planetary albedo drop. Gorodetskaya et al (2006) have used satellite sea ice cover and shortwave data to estimate the albedo drop for northern hemisphere sea ice regions. They obtain a 0.22 planetary albedo change for a 100% change in sea ice cover. This is roughly half the surface albedo difference between a typical sea ice cover and seawater.

While (1) is always nonlinear, the steepness of the ramp, the drop divided by the ramp temperature range, impacts the character of the nonlinearity. In particular, if the ramp is so steep that, as warming occurs, the extra shortwave absorption exceeds the extra loss of energy from OLR, the total feedback will be positive and there will be unstable transitions between ice-covered and ice-free states. This is an example of the slope-stability theorem of energy balance models (see Crowley and North, p. 18—19, for an elementary discussion). We can form an expression for the critical ramp temperature range, ΔT_C , between stable and unstable solutions:

$$\Delta T_C = S \Delta \alpha / B \quad (2)$$

A larger range is needed for stabilization when the insolation and the albedo drop are large and when the longwave damping is small.

When the ramp steepness is supercritical, the total feedback is positive in the ramp temperature range so it will contain only unstable equilibria. As a result, the ramp range becomes a forbidden zone, inaccessible with any forcing. As forcing is slowly varied, these temperatures are skipped leading to a discontinuity in polar temperature. Since the polar temperature contributes to the global mean temperature, it would also have a (much smaller) temperature discontinuity when the ramp steepness is supercritical.

If we insert the insolation at the North Pole (173 W/m^2), the Gorodetskaya et al planetary albedo drop, and a satellite-estimated outgoing longwave (OLR) damping value ($1.5 \text{ W/m}^2/\text{K}$, Marani 1999) into (2), we get a critical transition temperature range of about 25°C . This value indicates instability of the transition to a sea ice free climate because the annual mean temperature over the perennial sea ice today is about -18°C and the perennial ice-free zone just beyond the maximum ice edge is at about 0°C . Since ice in this temperature range is producing the satellite-observed planetary albedo drop, the average slope of the albedo/temperature curve on the way to ice-free conditions exceeds the critical slope and so the critical slope would have to be exceeded at some point, producing instability. However, this conclusion depends upon the assumption that heating from atmospheric transport remains fixed.

But it is unlikely that atmospheric heat transport would not respond to changes in shortwave absorption. The region north of 70°N receives more energy from the atmospheric transport than it absorbs from the sun and together they make up nearly all of the OLR – the surface flux is small (Serreze and Barry, 2005). Since atmospheric heat

transport is a big player in Arctic climate, it would not likely stand on the sidelines letting OLR completely balance a large change in absorbed shortwave. Indeed, it is probable that the Arctic heat convergence is as high as it is because it is countering the smallness of Arctic shortwave absorption, which is about 1/3 of the global mean (Serreze and Barry, 2005).

Horizontal temperature diffusion is a simple method of representing heat transport in an EBM. However, adding diffusive heat transport to the EBM does not eliminate unstable transitions in all cases. Instead, these diffusive transport models can exhibit an unstable loss of a finite patch of polar ice as forcing is increased. The instability is called the *small ice cap instability* (SICI) and in some ways is a companion to the large ice cap instability whereby the globe becomes ice-covered after the ice reaches a critical maximum extent. The ice edge lies in a temperature boundary zone having a length scale determined by the diffusivity and longwave damping parameters (North 1984). Both instabilities occur when this zone impinges on a boundary, either the equator or the pole. The instability can be removed by reducing the albedo ramp slope but the main point here is that the instability can occur in spite of down-gradient (warm to cold) transport.

Furthermore, we expect that the transport changes in response to CO₂ increase will have a significant up-gradient component. In the atmosphere this comes about because warmer air allows for an increase in the latent heat transport. Held and Soden (2006) show that increased latent transport drives an increase in heat transport to the polar regions, in spite of enhanced warming there, in both equilibrium and transient CO₂ increase experiments. Additionally, Holland and Bitz (2003) have shown that the ocean also transports more heat into the Arctic, even as the heat transport is being reduced at

lower latitudes in association with the weakened meridional overturning circulation. Thus, it is not clear that heat transport can be relied upon to stabilize Arctic climate by exerting a cooling influence on the region as it warms at a larger-than-global rate.

When evaluating the linearity of polar climate it will be useful to note the well-established fact that the global temperature response to forcing is linear. This was clearly shown for the GISS model by Hanson et al (2005) who calculated forcing efficacy, the ratio of global temperature change to forcing magnitudes for various forcing types and magnitudes. The efficacy was constant over a large range of magnitudes including the last glacial maximum and the anthropogenic future.

We can make use of the global linearity as follows, since global temperature change is linear in forcing, if polar temperature change is linearly related to it, then polar temperature change must also be linear. The ratio of polar to global temperature change is called the polar amplification. It is typically larger than one for a number of reasons including the ice-albedo feedback. If the polar amplification is also constant then polar climate change is linear. If there is a nonlinear relationship between polar and global temperature, a non-constant polar amplification, then the polar change must be nonlinear. Unstable behavior is a subcategory of nonlinear behavior. If the relationship between polar and global temperature is nonlinear *and* shows a temperature discontinuity, we have evidence of an unstable polar climate change.

III. Arctic linearity in 21st century experiments

Now we turn to the GCMs to see whether the projected 21st century polar climate change exhibits nonlinearity. Since Arctic climate is quite variable, it will be useful to do some averaging to bring out the forced signal. First, to form ΔT_P , the change in polar

surface air temperature, we average over the “half-cap” polar region north of 80N between 90E and 270E in the coldest part of the Arctic Ocean. In the remainder of this paper this region is referred to as the *polar* region. Next we take 5-year averages and we also average over the separate runs of the individual models made available in the PCMDI archive, thus each point represents an average over 15 to 35 years, depending upon the model’s ensemble size. Fig. 1 shows the results for five models that supplied multiple runs to the archive for the SRES A1B experiment. In spite of differences in global warming, polar warming and polar amplification, all of the models show a very linear relationship between polar and global temperature change. From this close relationship, it is clear that simulated 21st century polar climate change is very linear. While linear, the polar temperature changes are quite large in some of the models, approaching the magnitude of the abrupt warmings of the Dansgaard-Oeschger cycles in the glacial era in magnitude.

Figure 2 shows the relationship between polar region effective albedo and surface temperature. The effective albedo is the long-term ratio of surface up to surface down shortwave flux. Effective albedo can be shown to be the time averaged albedo weighted with the surface downward shortwave. This weighting is especially important in the Arctic where the insolation has a very large seasonal cycle. In general, the albedo and temperature have close linear relationships. Four of the five models become ice-free during September in the polar region over the course of the 21st century without disturbing this relationship. In terms of the EBM discussion of the last section, the simulated Arctic climate changes are linear because the albedo/temperature relationship is entirely within a subcritical linear ramp region in the 21st century.

The NCAR CCSM3 model, which spans the largest range of temperatures and albedos, shows some gentle downward arcing in its albedo/temperature relationship (Fig. 2). Since this arcing is not apparent in the polar amplification plot, other feedbacks must be compensating for its slightly nonlinear impact. An examination of a similar plot for the planetary albedo (not shown) does not show this arcing behavior, so the compensation may occur between atmospheric and surface shortwave terms.

Holland et al (2006) have noted that there are abrupt declines in September Arctic sea ice cover in the individual ensemble members of the NCAR CCSM3 SRES A1B experiments. Part of the steepness of the ice cover decline must be related to an acceleration of global warming in the early 21st century under SRES A1B forcing. Holland et al report that the annual mean ice cover in the NCAR CCSM3 is linearly related to global mean temperature, a result earlier found in the UKMO HadCM3 model by Gregory et al (2002), but that September ice cover is not so related. Therefore another factor must be involved in these sharp declines. They note that ice cover responds more sensitively to melting when it is thin. Part of the acceleration of the ice cover decline and its increase in variability are likely due to this increased sensitivity. It is to be expected that a binary variable such as ice cover will show some degree of nonlinearity when confined spatial and temporal averaging is done. The abrupt September ice cover declines are perhaps best characterized as a nonlinear response to linear climate dynamics.

IV. Arctic nonlinearity in annually sea ice-free experiments

From Fig. 2 we note that, even with the complete loss of September ice in most of the models, effective albedo has a long way yet to fall to approach open water values of

about 0.1. Furthermore, following their linear trends, the models would achieve this albedo at temperatures well above freezing – between 11°C and 29°C. The curves would therefore likely experience considerable steepening under further warming – potentially inducing nonlinear climate changes.

We can only be sure of observing the presence or absence of nonlinear behaviors associated with ice-albedo feedback in experiments that warm to the point of complete ice removal. Beyond this point there can be no further reductions in polar ocean surface albedo. The presence or absence of sea ice is easily determined by examining air temperatures in the coldest month and annual effective surface albedos (the ratio of annual surface-up to annual surface-down shortwave fluxes). If the coldest month temperature is at freezing and the effective albedo is near an open ocean value (about 0.1) then we can be assured that there is little sea ice in the particular region in either summer or winter. Seventy nine runs of four standard experiments (1%/year CO₂ increase to doubling, 1%/year CO₂ increase to quadrupling, SRES A1B and SRES A2) were examined for annually ice-free conditions in their polar regions (80N-90N, 90E-270E) based on these criteria. Of these, only two, had years with February polar region temperatures at freezing temperature and annual surface albedos below 0.15. Thus, it is quite uncommon for a model's Arctic ocean to become sea ice-free year around in these climate change experiments. By contrast, it is common in these runs for the Arctic sea ice to disappear in September – about half of the runs had Septembers with surface albedos less than 0.15. Unlike Thorndike's "toy" model, the seasonally ice-free state is apparently quite stable in GCMs.

The two runs which lose their Arctic sea ice year around are the 1%/year CO₂ increase to quadrupling experiments of the MPI ECHAM5 and the NCAR CCSM3.0. Eleven other models supplying data for this experiment did not lose all Arctic sea ice. Of the four forcing scenarios, the quadrupling experiment attains the highest forcing level, over 7 W/m². Both models are run for nearly 300 years, well past the time of quadrupling at year 140. The atmospheric CO₂ is held constant after quadrupling but temperatures are generally still rising in the models as the ocean heat uptake declines (Stouffer 2004).

The top panel of Figure 3 shows surface albedo for the polar region as a function of time for the two model experiments. The seasonal ice state is indicated by the albedos for three months: March (blue), June (green) and September (red). The annual effective albedo (light blue) characterizes the time-mean reflective capacity of the ice pack. The NCAR model loses its September sea ice near year 50, the MPI model loses it later, at about year 100. Both models have a progression of albedo reductions moving to earlier months in the sunlit season over the course of the integrations. The March sea ice is lost abruptly in the MPI model in the CO₂ stabilized period. The March decline is more gradual in the NCAR model. The variability of March albedos after the decline indicates occasional reappearance of ice in the NCAR model but not in the MPI model.

The lower panel of Figure 2 shows the albedos as a function of polar region surface air temperature. The albedo changes are more similar when viewed as a function of temperature rather than as a function of time but differences remain. The MPI model albedo declines are more abrupt in temperature as well as in time. Both models become seasonally ice-free (September albedo goes flat) at an annual polar temperature of about -

9°C. It is noteworthy that this loss of ice does not alter the nature of the decline in effective annual albedo in either model. This behavior was also noted in the 21st century plot (Fig. 2). Here we see that the extension to warmer temperatures does involve nonlinear effective albedo changes – steepened arcing in the CCSM3 and a kink-like turn in ECHAM5. The total fall in effective surface albedo is over 0.5 in both models but the effective planetary albedo drop over the experiments is about 0.1 for both models (not shown), indicating a large role for atmospheric shortwave masking and shortwave property changes.

The rapidity of the transition to annually ice-free conditions in the ECHAM5 model and the failure of subsequent variability to produce significant ice are suggestive of an unstable transition to a new equilibrium. Since the lifetime of sea ice in the Arctic (about 10 years) is short compared to the timescale of CO₂ increase (70 years for CO₂ doubling), we can view the ice as passing through a series of quasi-equilibrated states as the warming progresses. Under this interpretation the rapid transition to the annually ice-free state in the MPI model bears some resemblance to the SICI of simple energy balance models which occurs abruptly as a global forcing is gradually raised above a threshold value.

To explore further the connection between the transition and SICI, we look at the changes in surface albedo feedback (SAF) as the transition progresses. Using the fact that the model transitions are more similar in temperature than in time (Fig. 3), we evaluate the surface albedo feedback in three (annual mean) temperature eras: -15°C to -10°C (perennial to seasonal ice transition), -10°C to -5°C (seasonal ice), and -5°C to 0°C

(transition to ice-free). The ALL/CLR method of Winton (2005) is used to estimate the SAF (see Winton 2006 for details).

In both models, surface albedo feedback makes an increasing contribution to the decline in sea ice as air temperatures approach freezing (Fig. 4a). In the NCAR model the increase is gradual, consistent with the arcing decrease in effective annual surface albedo (Fig. 3). In the MPI model a sharp increase occurs in the transition to ice-free temperature range, consistent with the kinked shape of the effective annual albedo decline for that model. In the MPI model, the SAF becomes very large ($2.3 \text{ W/m}^2/^{\circ}\text{C}$) in the warmest temperature range.

Figure 4b shows the monthly contributions to the SAF of the two models in the three temperature ranges. As the warming progresses, there is a shift to earlier months in the sunlit season. This shift allows the SAF to increase even as the ice-free season appears and grows. Aside from seasonal insolation variation, the early months of the sunlit season potentially contribute more to SAF than the later months for two reasons:

(1) Surface albedos are initially larger so there is the potential for a larger albedo reduction as the ice is removed exposing the low albedo seawater. Figure 3 shows that September albedos are 0.1 to 0.2 lower than those in March at the beginnings of the runs.

(2) Atmospheric transmissivities are largest in the spring and decline through the summer to a minimum in September in both models.

Ignoring multiple cloud-ground reflection, the SAF is the product of the downward atmospheric transmissivity and the surface albedo change, so these two factors compound each other.

The pattern of surface albedo decline in the CCSM3.0 model (not shown) shows a plume of reduced albedo penetrating into the half-cap region from the Kara Sea, indicating an oceanic influence in the decline. This interpretation is borne out by an examination of the polar region heat budgets for the two models shown in Fig. 5. These budgets are constructed by regressing the fluxes on temperature in the three temperature ranges. The slopes are then multiplied by 5°C to give a representative flux change between the beginning and end of the specified temperature era. Figure 5 shows that the large increase of SAF in the MPI model at warmer temperatures also appears in the overall shortwave budget of the region and that there is a smaller increase in the shortwave budget of the NCAR model. The outgoing longwave radiation has a small damping effect on the warming of the region in both models. The surface budget changes are quite different: the MPI model has only small changes while the NCAR model has a large increase in surface forcing as the ocean supplies increased heating. This ocean heating contributes more to the warming of the NCAR model in the warmest temperature era than the SAF. The atmospheric heat transport convergence shifts from a forcing for the warming in the coldest temperature era to damping the warming in the two warmer eras in both models. It is this change in atmospheric convergence of heat, rather than the OLR, that does the most to balance the forcing factors: shortwave in MPI, shortwave plus surface in NCAR. All of the surface flux changes are opposite to the atmosphere flux convergences in their impacts on the warming.

To gain a sense of the regional extent of nonlinear climate changes, we split the polar amplification into two factors: polar to Arctic (60N-90N) and Arctic to global amplifications. Fig. 6 shows the relationship between polar and Arctic temperatures (top)

and Arctic and global temperatures (bottom) for the two models over the course of the 1%/year to 4X CO₂ runs. The warmest polar temperature attained in the two models is about the same but the global temperature rise is considerably larger in the MPI model, while the Arctic/global and polar/Arctic amplifications are correspondingly smaller. The lines in Fig. 6 are fits to the relationships for data with polar temperatures less than -5°C. A deviation from this fit at warmer temperatures might reflect the enhanced warming due to the dramatic changes in sea ice cover above this temperature in both models. The relationship is mainly linear in both models but in the ECHAM5 model the polar temperature rises above the reference line starting at a polar temperature of -4°C until it is about 2°C larger and then begins to parallel the fitted line at a polar temperature of 0°C. Apparently the large increase in surface albedo feedback in this range of temperatures (discussed above) plays a role in this extra warming of the polar region. After the ice is eliminated, the SAF drops to zero, and further warming falls below the -4°C to 0°C ratio. The behavior of the CCSM3.0 is somewhat different. At -5°C the polar temperature rises slightly above the fitted line but then parallels it as both regions warm further. In both cases, the transition to seasonally ice free at a polar temperature of -9°C does not disturb the linear relationship between warming in the two regions. The relationship between Arctic and global temperatures (Fig. 6, bottom) is quite linear in both models indicating that the nonlinear changes in the Arctic ocean do not have significant impacts on the broader region temperatures. Although the elimination of Arctic sea ice would doubtless have enormous consequences for the local environment, these models do not show it to be particularly important for the larger scale climate changes.

V. EBM interpretation of the transition to annually ice-free

Now, to provide a mechanistic comparison to the GCM behaviors of the last section, we examine polar amplification in a simple one dimensional EBM as it experiences small ice cap instability. Following North (1984), the temperature equation for the EBM is:

$$-D \frac{d}{dx} (1-x^2) \frac{dT}{dx} + A + BT = S(x)[1 - \alpha(T)] \quad (3)$$

Table 1 defines the notation and gives parameter values. The value for the longwave sensitivity parameter, B , comes from a regression of ISCCP outgoing longwave on surface temperature (Marani 1999). The value used for the albedo jump with temperature is the Gorodetskaya et al (2006) value for the radiative effectiveness of northern hemisphere sea ice – the impact on planetary albedo of the change from total to zero sea ice cover. This was determined using ERBE shortwave measurements and HadISST1 sea ice concentration data. After setting these two parameters, the atmospheric diffusivity, D , is adjusted to give a reasonable planetary range of surface temperatures.

Most EBM studies explore climate sensitivity by varying the solar constant. Here, we are interested in exploring the relationship between temperature change at the pole ($T(x=1)$) and the global mean temperature change ($\int_0^1 T(x)dx$) as climate warms. To this end, it is desirable to force in a manner that does not affect this relationship. So here we force climate in a meridionally uniform way by varying A – reducing A induces warming. Thus, we can think of A as a forcing for global mean temperature since

$$\int_0^1 T(x)dx = \left(\int_0^1 (1-\alpha)Sdx - A \right) / B.$$

With this forcing, all polar amplification is due to ice-albedo feedback, the only positive feedback in the system.

Initially we configure the EBM with a step jump in albedo at 0°C ($\Delta T_M=0^{\circ}\text{C}$). North used -10°C as the location of the step change. This lower value presumably represents the temperature needed to retain terrestrial snow through the summertime. Sea ice has a source (seawater freezing) that decreases with increased temperature but is positive while there are periods of below freezing temperatures. This added source is a factor aiding the persistence of summer sea ice at higher annual temperatures than terrestrial snow.

Fig. 7 shows the polar and global temperatures for the MPI ECHAM5 experiment discussed in the previous section (green) and the EBM with a step albedo jump (blue) and with the same jump smoothed over a transition zone of 5°C (red). The EBM changes have been forced by varying A in (3) while the GCM changes are forced by CO_2 increase, of course. The CO_2 forcing itself is generally somewhat reduced in the Arctic (Winton 2006a). Nonetheless, the GCM line is the steepest at each polar temperature, so the polar amplification is always larger for the GCM than for the EBM. This is consistent with the finding of a number of studies that factors beside the surface albedo feedback contribute significantly to polar amplification of climate change (Alexeev 2003; Holland and Bitz 2003; Hall 2004; Winton 2006a). Further evidence of this can be seen in the MPI ECHAM5 curve where significant polar amplification remains even after the sea ice has been eliminated. The EBM does not represent these additional factors and so has smaller polar amplifications.

The EBM with a step albedo change has a discontinuity in polar and global temperatures where the small ice cap instability is encountered and both warm abruptly with the removal of the reflective ice cap. The light blue dashed line spans this jump and its slope defines a polar amplification across the instability. In the cooler part of the

curve, to the left of this jump, the pole is always below freezing temperature so the local shortwave absorption does not change. The amplification of polar temperature change over global in this part of the curve, about 1.8, is due to the influence of increased absorption of shortwave energy at the ice edge, as the ice retreats poleward, conveyed to the pole by atmospheric transport. North (1984) shows that, as the instability is approached, the pole feels nearly as much warming impact from the ice retreat as the ice edge itself.

Based on the fact that the ice cap covers about 6% of the hemisphere before its elimination, we might expect the polar amplification across the jump to be about 16, since this increased absorption is the cause of both temperature jumps. The actual polar amplification is much less because atmospheric heating at the pole, which has been increasing to that point, collapses with the ice cap, countering its local impact to a large degree (Fig. 8). After the ice cap collapse, there is no ice-albedo feedback and polar and global temperatures rise in a one-to-one relationship. The sequence of changes in polar energy budget encountered as the climate warms lead to a medium/high/none sequence of polar amplifications in the EBM.

The global and polar temperatures for the MPI ECHAM5 show a three slope regime behavior similar to that of the EBM. However, the GCM does not show any discontinuity in these temperatures. This may be partly due to the GCM, unlike the EBM, not being fully equilibrated at each point in time and hence able to fill the “forbidden zone” with transient temperatures. However, taking note that the ECHAM5 sensitivity of polar albedo to temperature is steep but far from step-like (Fig. 3b), we explore the possibility that having the albedo changes occur over a finite range of

temperatures stabilizes the transition while retaining enhanced polar sensitivity due to increased surface albedo feedback. EBM runs show that the multiple equilibria remains, with a reduced ΔT across the jump, for a ramp range of $\Delta T_M = 4^\circ\text{C}$ but is eliminated when for $\Delta T_M = 5^\circ\text{C}$. The plot for the polar amplification in the stabilized case (Fig 7, red) shows that a continuous section of enhanced polar amplification fills the region occupied by the jump in the step albedo EBM. The enhanced sensitivity in this region is caused by the reduced overall (negative) feedback due to a positive, but subcritical, local ice-albedo feedback. Fig. 8 shows that the diffusive term, operating as a negative feedback, provides less heating to the pole, opposing the enhanced shortwave absorption.

The change of atmospheric heat transport convergence with polar temperature in Fig. 7 shows that, in the GCMs as well as the EBM, enhancement of polar warming by atmospheric transport at low temperatures gives way to a damping impact at higher temperatures. However, the GCM transition occurs at much lower temperature, perhaps partly due to its having a polar albedo response at lower temperatures. Other mechanisms, not present in the EBM, can significantly impact poleward transport in the GCMs, for example the enhancement of latent heat transport with temperature (Alexeev 2003; Held and Soden, 2006). As the ice-free state is approached, the damping effect of the atmospheric heat transport change is much larger than longwave damping in the EBM as in the GCMs.

VI. The tethering effect of heat transport

The previous section shows that atmospheric heat transport plays an important but complicated role in polar climate change – initially forcing the region to warm at a greater-than-global rate but eventually becoming a cooling influence at higher

temperatures. Held and Soden (2006) show that the latent heat component of the transport scales up in a warming climate according, roughly, to the Clausius-Clapeyron relationship. This increase drives an increase of the total transport toward the North Pole in spite of polar amplification. However, it is possible that, even in the early warming, part of the transport is helping to maintain the very constant polar amplifications seen in Fig. 1. To reveal this moderating role we perform two diagnostic experiments that force only the polar regions and examine the damping mechanisms.

The first is a modification of the AMIP experiment – an atmospheric GCM with specified sea surface temperatures and sea ice cover. We perform a twin to this experiment where the sea ice boundary condition is replaced with seawater freezing temperature and albedo. The experiment is done with the atmospheric component of the GFDL CM2.1 climate model. A similar experiment for the DJF season was performed earlier by Royer et al (1990). The impact on atmospheric temperatures and winds in the current experiment are in general agreement with those found by Royer et al. Figure 9 shows that there is an intense warming of the lower polar atmosphere, mainly confined below the 0°C potential temperature contour of the control, “ice-in”, simulation. This is consistent with the regionally limited response of the GCMs to transition to ice-free conditions shown in Fig. 6. Other features found by Royer et al, an equatorward shift of the jet, redistribution of sea level high pressure away from the central Arctic to adjacent land regions, and a reduction of cloud cover as the Arctic Ocean becomes more convective, are also found in this experiment (not shown).

Our main interest in the experiment is to assess the stability of the Arctic ice and its causes. The net surface heat flux change in the ice-covered regions is the result of two

competing changes: (1) increased shortwave absorption due to lowered albedo, and (2) increased longwave and turbulent heat loss due to increased surface temperature. The net upward heat flux change of 24 W/m^2 (Fig. 10) indicates that (2) is dominant so the ice is stable and would grow back at an initial rate of 2.5 m/yr . At the top of the atmosphere, the extra shortwave absorption is only partly balanced by increased OLR. Most of the damping influence comes from a reduction of heat transport into the ice covered region by the atmosphere. This reduced heat transport, in turn, is mainly supported by a reduction in surface heat flux from the adjacent ice-free ocean, particularly in the North Atlantic (Fig. 11).

The AMIP experiment fixes SSTs implicitly assuming that the ocean has an infinite heat capacity. This assumption may be reasonable here since the near-ice regions that are experiencing large heat flux changes are occupied by deep wintertime mixed layers with much larger heat capacity than the sea ice or the shallow atmospheric layer that interacts with it. Nonetheless, it is useful to relax this assumption by performing a similar experiment in a fully coupled climate model – a developmental version of GFDL CM2.1. In this 100 year experiment, we force the ice region by lowering the ice albedos. Fig. 12 shows changes in climate model heat fluxes as in Fig. 10. As expected, there is an increase in shortwave absorption at the top of the atmosphere and, as in the AMIP case, it is only partially offset by a local OLR change. Again the main balancing effect is from a reduction of atmospheric heat transport convergence into the ice-covered region (defined from the control experiment). There is also some net downward flux at the surface in the ice-covered region which is mainly supported by reduced latent heating due to the reduced sea ice export in thinner ice. This change in sea ice transport has a salinifying

influence near the sea ice edge. Again, the reduced heat transport into the Arctic is supported by a reduction in heat flux out of the adjacent ocean surface. But this response is now oversized compared with the changes in the sea ice region. The reason for this is that, in spite of the change in sea ice freshwater forcing, reduced ocean heat extraction caused by reducing sea ice albedo has induced a reduction of the meridional overturning circulation (MOC). This is shown in Fig. 13 along with the change in deepwater ages averaged over the 100 years. The reduction in deepwater ventilation agrees with the result of a similar sea ice albedo reduction experiment performed by Bitz et al (2006) with the NCAR CCSM3. The CCSM3 response was relatively larger in the southern ocean, perhaps because the CCSM3 has more southern ocean sea ice to feel the albedo reduction than GFDL's CM2.1.

VII. Summary and discussion

The potential for sea ice-albedo feedback to give rise to nonlinear climate change in the Arctic ocean – defined as a nonlinear relationship between polar and global temperature change or, equivalently, a time-varying polar amplification has been explored in the IPCC AR4 climate models. Five models supplying SRES A1B ensembles for the 21st century were examined and very linear relationships were found between polar and global temperatures indicating linear Arctic climate change. The relationship between polar temperature and albedo is also linear in spite of the appearance of ice-free Septembers in four of the five models.

Two of the IPCC climate models have Arctic Ocean simulations that become annually sea ice-free under the stronger CO₂ increase to quadrupling forcing. Both runs show increases in polar amplification at polar temperatures above -5°C and one exhibits

heat budget changes that are consistent with the *small ice cap instability* of simple energy balance models. Both models show linear warming up to a polar temperature of -5°C , well above the disappearance of their September ice covers at about -9°C . Below -5°C , surface albedo decreases smoothly as reductions move, progressively, to earlier parts of the sunlit period. Atmospheric heat transport exerts a strong cooling influence during the transition to annually ice-free conditions.

Specialized experiments with atmosphere and coupled models show that perturbations to the sea ice region climate are opposed by changes in the heat flux through the adjacent ice-free oceans conveyed by altered atmospheric heat transport into the sea ice region. This, rather than OLR, is the main damping mechanism of sea ice region surface temperature. This strong damping along with the weakness of the surface albedo feedback during the emergence of an ice-free period late in the sunlit season are the main reasons for the linearity of Arctic climate change and the stability of seasonal ice covers found in the IPCC models.

Support for these mechanisms can be found in simple models. Thorndike's "toy" model shows their importance by demonstrating the consequences of their absence. In order to make the "toy" model analytically tractable, the ice experiences a constant insolation in the summer season. This disables the seasonal adjustment of albedo feedback just mentioned. Furthermore, the Thorndike model has only weak surface temperature damping to space through a gray-body atmosphere. Atmospheric heat transport convergence is held fixed. The result is a model that has no stable seasonal cycle with ice-covered and ice-free periods -- the model is either annually ice-covered or annually ice-free. Eisenman (2006) has enhanced the Thorndike model with a sinusoidal

insolation and temperature sensitive atmospheric heat transport and finds stable seasonally ice-free seasonal cycles over a broad range of CO₂ forcings, from 2 to 15 times current levels.

The stabilization of sea ice through opposing heating or cooling of the surrounding ocean was shown to have substantial impact on the overturning circulation and deep ventilation. This effect may provide a partial answer to an outstanding question raised by the CMIP group experiments (Gregory et al, 2002). These experiments separate MOC weakening under CO₂ increase into freshwater and thermally induced components by using two auxiliary experiments: control radiation with CO₂ increase ocean freshwater forcing and the reverse. These experiments show that the two effects basically add linearly and that thermal forcing dominates the MOC response. The mechanisms for this thermal impact are yet to be explored but ice-albedo feedback, and polar amplification more generally, may be one of them.

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Captions

Figure 1: Polar vs. global temperature for GCM ensembles forced with SRES A1B scenario. Each point represents an average over a 5-year period and over all ensemble members.

Figure 2: Polar effective albedo (annual surface up to surface down ratio) vs. polar temperature for GCM ensembles forced with SRES A1B scenario. Each point represents an average over a 5-year period and over all ensemble members.

Figure 3: Polar region albedo as a function of time (top) and annual mean polar region surface temperature (bottom) for the MPI ECHAM5 (circles) and NCAR CCSM3.0 (plusses) models. All data have been 5-year boxcar filtered.

Figure 4: Polar surface albedo feedback in three temperature eras (top). Monthly contribution to polar surface albedo feedback (bottom) for surface temperatures less than -5C (dashed) and between -5 and 0C (solid) for the NCAR CCSM3.0 (black) and MPI ECHAM5 (gray) models. All data have been 5-year boxcar filtered.

Figure 5: Polar atmosphere heat balance changes over three temperature eras: top-of-atmosphere absorbed shortwave (upper-left) and outgoing longwave radiation (upper right), atmospheric heating from sides (lower-left) and heating from the surface (lower-right).

Figure 6: Polar vs. Arctic temperature (top) and Arctic vs. global temperature (bottom) for MPI ECHAM5 (circles) and NCAR CCSM3.0 (plusses). All data have been 5-year boxcar filtered.

Figure 7: Polar vs. global temperature for MPI ECHAM5 (green circles) and the EBM (step albedo, – blue line; smoothed albedo – red line). ECHAM5 data have been 5-year boxcar filtered.

Figure 8: Polar atmospheric heat transport convergence plotted against polar temperature for EBMs with step albedo (blue) and smoothed albedo (red).

Figure 9: Change in zonal mean temperature due to sea ice removal in an AMIP experiment with the GFDL AM2.1 model. The freezing potential temperature contour of the “ice-in” experiment is shown for reference.

Figure 10: Integrated heat flux changes for the northern hemisphere ice-covered and ice-free regions (defined by the “ice-in” annual mean ice concentration).

Figure 11: Change in heat fluxes at the top (a) and bottom (b) of the atmosphere and the change in atmospheric heat transport convergence (c) for the AMIP ice removal experiment.

Figure 12: As in Fig. 10 but for the coupled model (GFDL CM2.1) reduced albedo experiment.

Figure 13: Change in coupled model overturning (contours) and 100 year mean age (shading) due to reducing sea ice albedo.

Table 1

x	sin(latitude)
T(x)	Surface temperature (°C)
A, B	Longwave parameters ($B=1.5 \text{ Wm}^{-2}\text{°C}^{-1}$; A is variable (Wm^{-2}))
D	Atmospheric diffusion ($=0.36B$, Wm^{-2})
$\alpha(T)$	albedo ($=0.3$ if $T>0^\circ\text{C}$; $=0.52$ if $T<-\Delta T_M$; $=(0.3(T+\Delta T_M)-0.52T)/\Delta T_M$, otherwise)
ΔT_M	Temperature range over which albedos transition linearly from ice-covered to ice-free
S(x)	annual shortwave distribution ($=340(1-0.482P_2(x))$; $P_2(x)=(3x^2-1)/2$)

Figure 1: Polar vs. global temperature for GCM ensembles forced with SRES A1B scenario. Each point represents an average over a 5-year period and over all ensemble members.

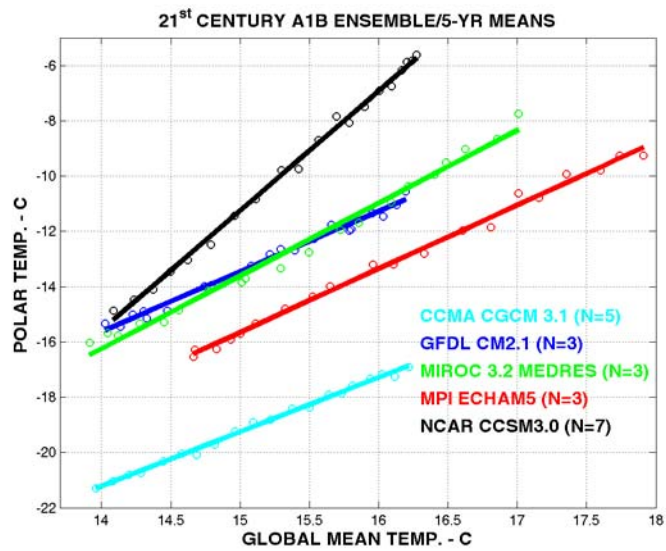


Figure 2: Polar effective albedo (annual surface up to surface down ratio) vs. polar temperature for GCM ensembles forced with SRES A1B scenario. Each point represents an average over a 5-year period and over all ensemble members.

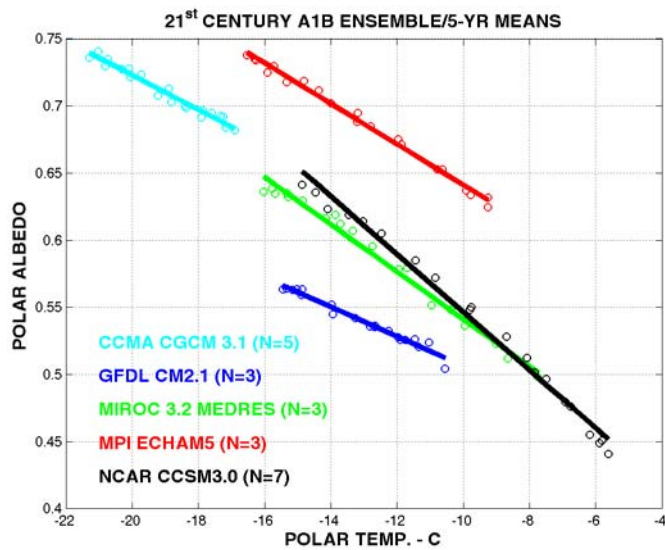


Figure 3: Polar region albedo as a function of time (top) and annual mean polar region surface temperature (bottom) for the MPI ECHAM5 (circles) and NCAR CCSM3.0 (plusses) models. All data have been 5-year boxcar filtered.

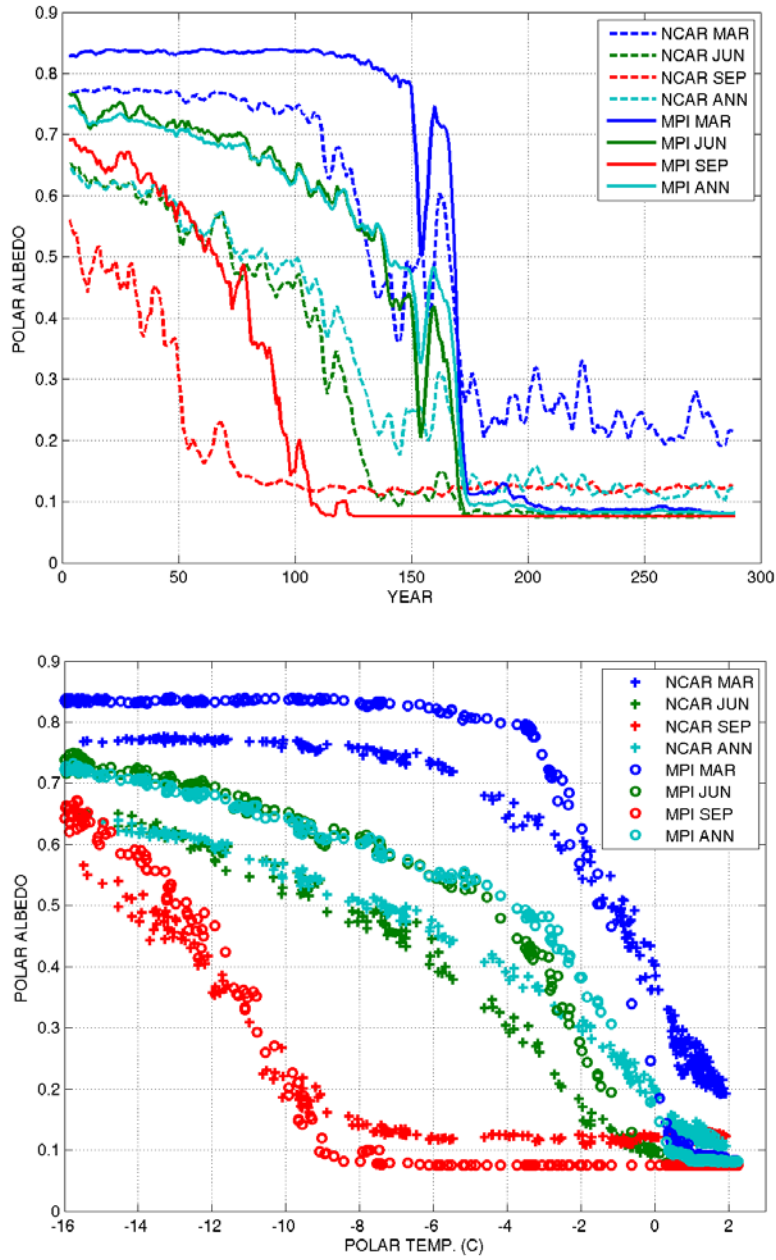


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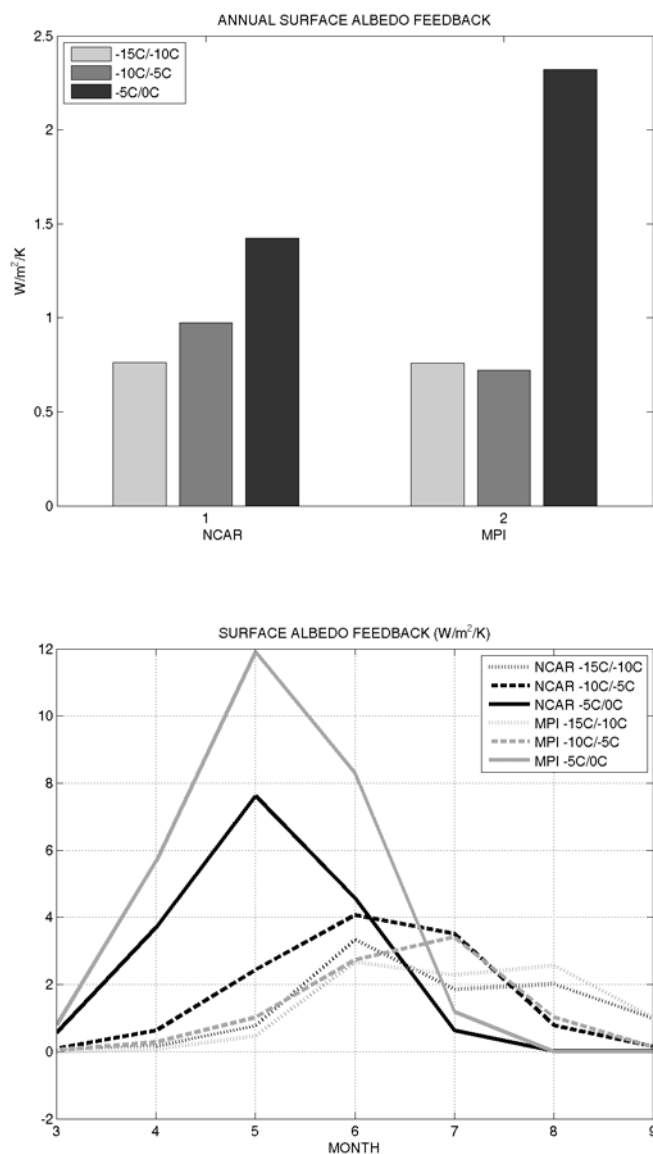


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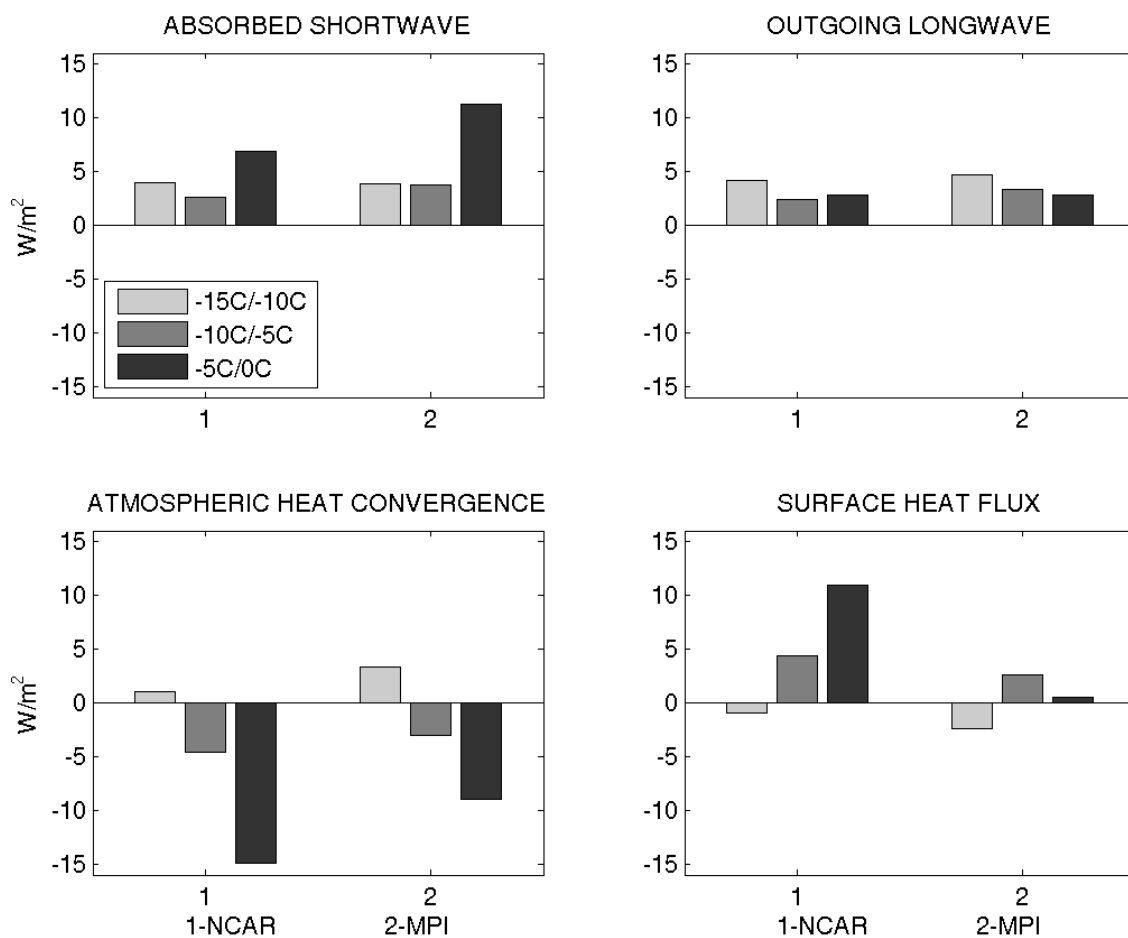


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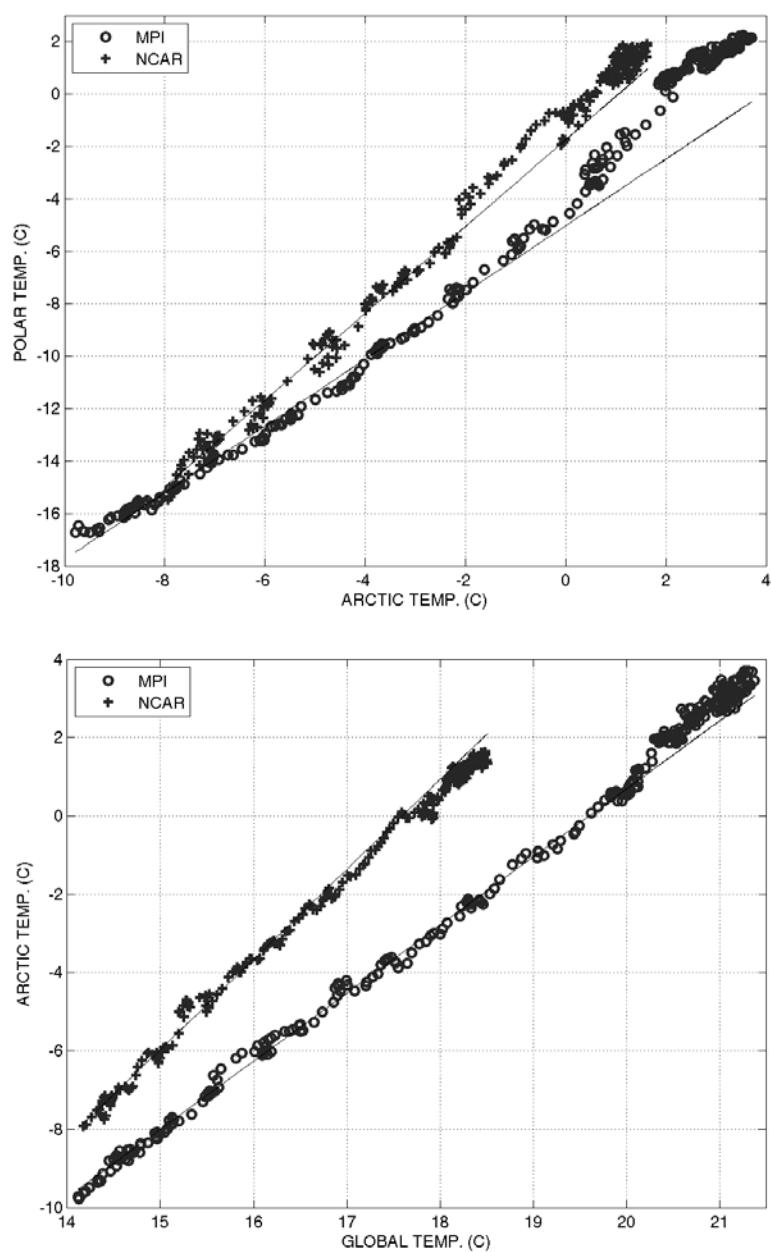


Figure 7: Polar vs. global temperature for MPI ECHAM5 (green circles) and the EBM (step albedo, – blue line; smoothed albedo – red line). ECHAM5 data have been 5-year boxcar filtered.

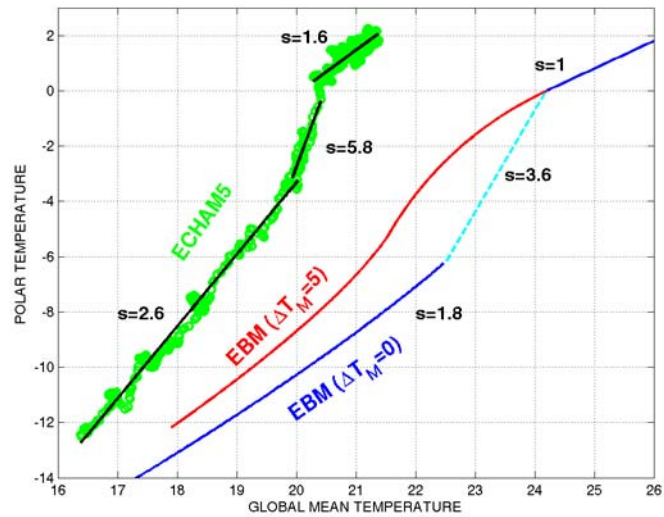


Figure 8: Polar atmospheric heat transport convergence plotted against polar temperature for EBM models with step albedo (blue) and smoothed albedo (red).

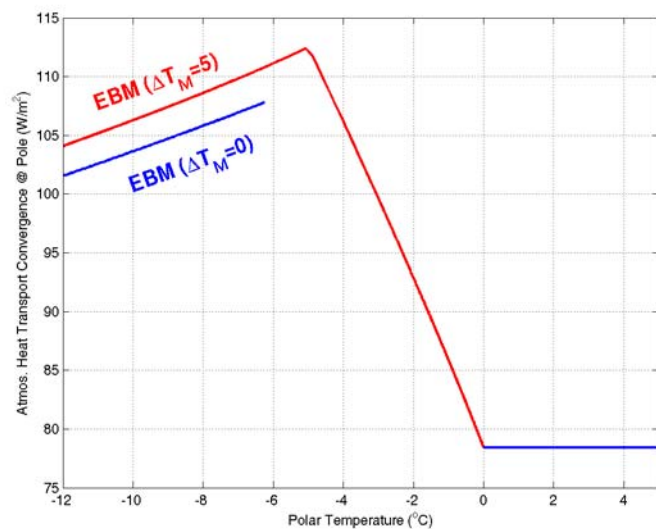


Figure 9: Change in zonal mean temperature due to sea ice removal in an AMIP experiment with the GFDL AM2.1 model. The freezing potential temperature contour of the “ice-in” experiment is shown for reference.

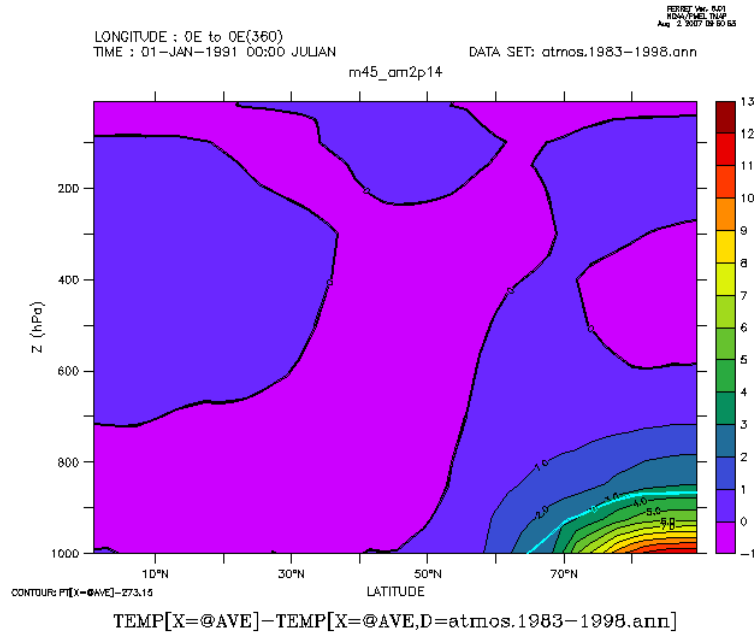


Figure 10: Difference in atmospheric heat fluxes over ice-covered and ice-free regions of the northern hemisphere between “ice-out” and “ice-in” AMIP runs. The “ice-covered” region is defined by the annual mean ice concentration of the “ice-in” experiment. All fluxes are in units of W/m^2 -ice-covered-region.

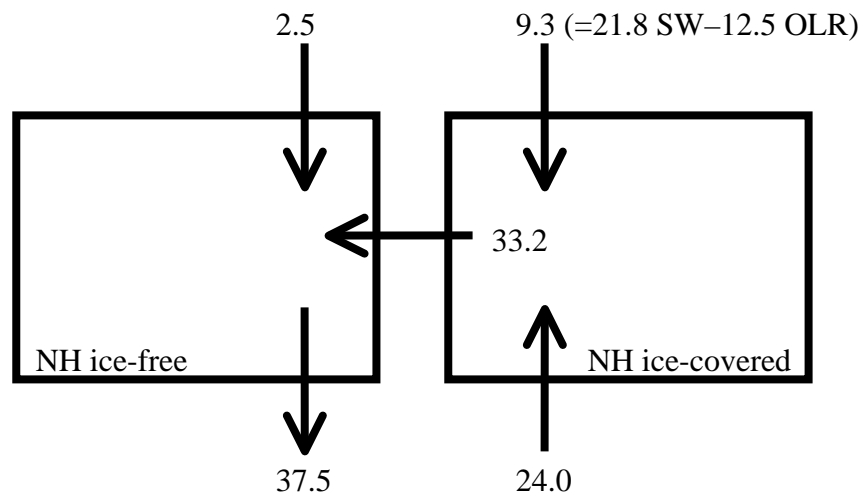


Figure 11: Change in heat fluxes at the top (a) and bottom (b) of the atmosphere and the change in atmospheric heat transport convergence (c) for the AMIP ice removal experiment. Panel (d) shows integrated heat flux changes for the northern hemisphere ice-covered and ice-free regions (defined by the “ice-in” annual mean ice concentration).

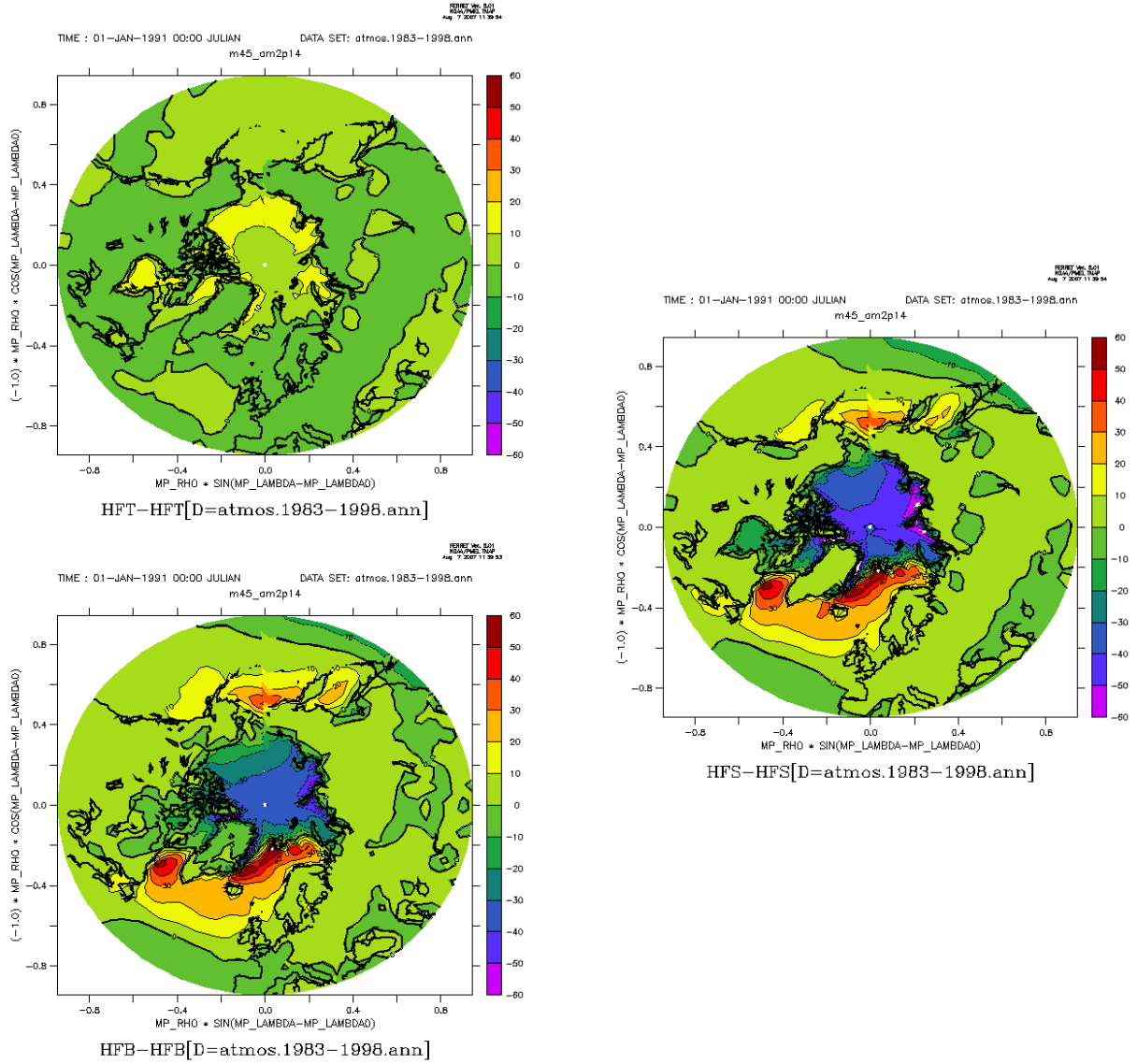


Figure 12: Difference in atmospheric heat fluxes over ice-covered and ice-free regions of the northern hemisphere between ice-albedo reduced and control coupled model runs. The “ice-covered” region is defined by the annual mean ice concentration of the control experiment. All fluxes are in units of W/m^2 -ice-covered-region.

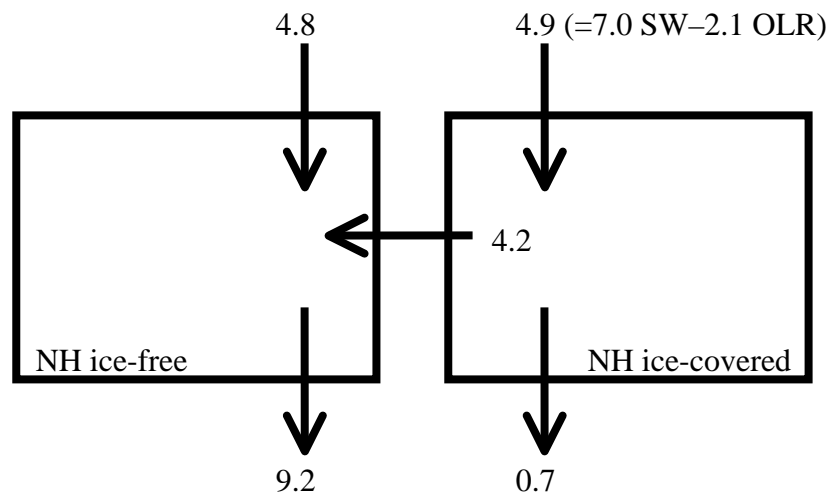


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